

Particle Collisions in Turbulence **The Role of Structures in the Flow**

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Extreme Dissipation and Intermittency in Turbulence

EUROMECH COLLOQUIUM (Virtual)

*Suspensions of sma*l*, heavy par*t*cles in a turbulent* fl*ow is commonplace*

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- $\dot{\mathbf{g}}$ \overline{c} \mathbf{r} Two examples is the third this function. $\frac{1}{2}$ Important questions relate to collisions and coalescences
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*Suspensions of sma*l*, heavy par*t*cles in a turbulent* fl*ow is commonplace*

The Carrier Turbulent Flow

 $\partial \mathbf{u}$ $\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} = \nu \nabla^2 \mathbf{u} - \nabla P + \mathbf{f}$

 $\nabla \cdot \mathbf{u} = 0$ **Incompressibility**

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The Carrier Turbulent Flow The Particle

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 $\nabla \cdot \mathbf{u} = 0$ **Incompressibility**

ARTICLES

Equation of motion for a small rigid sphere in a nonuniform flow

Martin R. Maxey^{a)} Department of Chemical Engineering, The Johns Hopkins University, Baltimore, Maryland 21218 James J. Riley

Flow Research Company, Kent, Washington 98031

(Received 11 June 1982; accepted 22 November 1982)

Maxey and Riley, Phys. Fluids (1983)

Basset, *Treatise on Hydrodynamics (1888)* Boussinesq. *Theorie Analytique de la Chaleur (1903)* Oseem*, Hydrodynamik* (1927) Tchen, PhD thesis, Delft (1947) Corrsin and Lumley, Appl. Sci. Res. (1956)

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$$
\frac{d\mathbf{x}}{dt} = \mathbf{v}
$$

$$
\frac{d\mathbf{v}}{dt} = \mathbf{g} - \frac{1}{\tau_p}(\mathbf{v} - \mathbf{u})
$$

The linear Stokes drag model

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$$

1100 100 110 Saw, Bewley, Bodenschatz, Ray, and Bec, Phys. Fluids Lett. (2016)

Working Equations

The Carrier Turbulent Flow The Particle

$$
\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} = \nu \nabla^2 \mathbf{u} - \nabla P + \mathbf{f}
$$

$$
\nabla \cdot \mathbf{u} = 0
$$

Characteristic timescale :
$$
\tau_{\eta} = \sqrt{\frac{\nu}{\varepsilon}}
$$

Characteristic lengthscale : $\eta = \left(\frac{\nu^3}{\varepsilon}\right)^{1/4}$

*d*v *dt* $=$ $\mathbf{g} - \frac{1}{\tau}$ τ_p $(\mathbf{v} - \mathbf{u})$ $d\mathbf{x}$ *dt* $=$ \bf{v}

 $\text{Characteristic timescale}: \tau_p =$ $2a^2\rho_p$ $9\nu\rho_f$

Characteristic lengthscale : *a*

Non-Dimensional Number $St =$ τ_p τ_η

Fingerprints

Turbulent Flows

- **• Irreversible and intermittent**
- **• Multiple length and time scales**
- **• Exotic structure: Vortices and straining regions**

Celani, *et al.***, Phys. Fluids (2001) M. Gupta and R. Singh**

Fingerprints

Turbulent Flows

- **• Irreversible and intermittent**
- **• Multiple length and time scales**
- **• Exotic structure: Vortices and straining regions**

Iner

- **Dissipative dynamics: Phase space contracts**
- **•** Spatial distribution strongly inhomogeners and the Concentration of Indian Increase o
	- **Maximum clustering achieved for finite**
	- **Uniform distribution in the small and lar** St=0.57

A. M. Wood, *et al.*, Int. J. Multiphase Flow, 31 (2005). E. Calzavarini, *et al.*, Phys. Rev. Lett., 101 (2008).

Longmire and Eaton, J. Fluid Mech. (1992) Bec, Phys. Fluids (2003) Bec, J. Fluid Mech (2005) Bec, Celani, Cencini, and Musacchio Phys. Fluids (2005)

E. Chun, *et al.*, J. Fluid Mech. (2005). Bec, *et al.*, Phys. Rev. Lett. (2007) Monchaux, Bourgoin, and Cartellier, Int. J. Multiphase Flows, (2012) Gustavsson and Mehlig, Adv. Phys. (2016)

The rate of collisions depends on the relative velocity of particles at contact

Saw, Bewley, Bodenschatz, Ray, and Bec, Phys. Fluids 26, 111702 (2014) James & Ray, Sci. Rep. 7, 12231 (2017) Bhatnagar, Gustavsson, and Mitra, Phys. Rev. E 97, 023105 (2018) Bhatnagar, Gustavsson, Mehlig, and Mitra,Phys. Rev. E 98, 063107 (2018)

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Tracers

- Determined by fluid velocity gradients: Increases with increasing Reynolds number
- Disregards whether the local velocity gradient arises from rotation or strain

Saffman & Turner, J. Fluid Mech. 1, 16 (1956)

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Inertial Particles

- Preferentially concentrate: Increases local particle density
- Sling effect: Large relative velocities through caustics
- Ideas as true for smooth random flows as it is for turbulent flows

Falkovich, Fouxon, and Stepanov, Nature 419, 151 (2002) Bec, Celani, Cencini, and Musacchio, Phys. Fluids 17, 073301 (2005) Wilkinson, Mehlig, and Bezuglyy, Phys. Rev. Lett. 97, 048501 (2006) Wilkinson, Mehlig, Östlund, and Duncan, Phys Fluids 19, 113303 (2007) Bec, S. Musacchio, and Ray, Phys. Rev. E 87, 063013 (2013) Gustavsson & Mehlig, Phys. Rev. E 87, 023016 (2013) Ravichandran & Govindarajan, Phys. Fluids 27, 033305 (2015) James & Ray, Sci. Rep. 7, 12231 (2017)

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These effects do not necessarily imply that collisions sense the structures of turbulence

PHYSICAL REVIEW E 93, 031102(R) (2016)

⊤erated coarescence

 2av^2 Ev $\frac{1}{2}$ λ∞ *ⁱ*−*j,j ni*−*^j nj* [−]! $\frac{1}{2}$ λ∞ *i,j ni nj .* (4) ² International Centre for Theoretical Sciences, Tata Institute of Fundamental Research, Bangalore 560012, India *SPEC, CEA Saclay, CNRS, 91191 Gif-sur-Yvette, France* (Received 8 July 2015; published 21 March 2016)

on particle inertia was actually devoted to estimating their is here shown to fail when the coalescing species are dilute and transport by a turbulent flow. The intermittent Lagrangian motion involves correlated violent events that lead to an unexpected rapid occurrence of the largest particles. This new phenomena is here quantified in terms of the anomalous scaling of turbulent three-point motion, leading to significant corrections in macroscopic processes that are critically sensitive to the early-stage emergence of large embryonic aggregates, as in planet formation or rain precipitation. balance. We thus get *n*˙2 ≃λ[∞]

For the next size, we have *n*˙3 ≃λ[∞] ¹*,*²*n*1*n*² and thus *n*3(*t*) ≃ *µo*lution of fluid elen dimensional turbulent flow during one large-eddy turnover time. Long Lagrangian evolution of fluid elements: Scaling and geometry tied together

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dominant destruction term [∝]*tⁱ*−1.

the velocity field and the associated nonuniqueness of fluid ele-

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The main assumption leading to Smooth Christian Library assumption assumption of the smooth contract of the smo
The main assumption of the smooth contract of the smooth contract of the smooth contract of the smooth contrac ition between the structure of the flow and droplet collisions and coalescences \vert λ[∞] *i,j* much faster than the evolution of *ni*. This is ensured for in one. Evidence suggests that how structure The implied correlation between the structure of the flow and droplet collisions and coalescences **|** remained an open one: Evidence suggests that flow structures matter, but how and when? $\qquad \qquad \blacksquare$ |*x*1(*t*) − *x*2(*t*)| ² [∼] ^ε*^t* 3. Still, when interested in more than two

What is the effect of structures in a flow on collisions and coalescence of **droplets?**

The Question

What is the effect of structures in a flow on collisions and coalescence of **droplets?**

Jason R. Picardo Lokahith Agasthya Rama Govindarajan PHYSICAL REVIEW FLUIDS **4**, 032601(R) (2019)

Rapid Communications

Flow structures govern particle collisions in turbulence

Jason R. Picardo,^{1,*} Lokahith Agasthya,^{2,†} Rama Govindarajan,^{1,‡} and Samriddhi Sankar Ray^{1,§} ¹*International Centre for Theoretical Sciences, Tata Institute of Fundamental Research, Bangalore 560089, India* ²*Indian Institute for Science Education and Research, Pune 411008, India*

The Answer

Straining regions more effective at generating collisions

PHYSICAL REVIEW FLUIDS **4**, 032601(R) (2019)

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The Answer

- Straining regions more effective at generating collisions
- Intense vorticity and strain, cohabiting as vortex-strain worm-rolls, conspire to generate violent collisions.

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The Answer

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Key Message: These effects are not *just* **due to preferential concentration**

PHYSICAL REVIEW FLUIDS **4**, 032601(R) (2019)

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Velocity gradient tensor

Symmetric: Strain Antisymmetric: Rotation

Direct Numerical Simulations

Velocity gradient tensor

Dubief & Delcayre, J. Turbul. 1, N11 (2000) Blackburn, Mansour, and Cantwell, J. Fluid Mech. 310, 269 (1996) Chong, Perry, and Cantwell, Phys. Fluids A 2, 765 (1990)

Direct Numerical Simulations

Velocity gradient tensor

Direct Numerical Simulations

Ficardo, Agasthya, Govindarajan, and Ray, Phys. Rev. Fluids (Rapids) 4, 032601(R) (2019)

or *rear-end* collisions, as opposed to *side-on* collisions, which are predominant in vortical regions. **Direct Numerical Simulations of the velocity gradient into the velocity gradient int**

Velocity gradient tensor

 $Q > 0$ [Vorticity] $Q < 0$ [Strain]

These can be measured along the particle trajectories Ireland, Bragg, and Collins, J. Fluid Mech. 796, 617 (2016)

Ficardo, Agasthya, Govindarajan, and Ray, Phys. Rev. Fluids (Rapids) 4, 032601(R) (2019)

Direct Numerical Simulations **Consequent** in Section of the velocity gradient into the velocity gradient into the velocity gradient in straining α or *rear-end* collisions, as opposed to *side-on* collisions, which are predominant in vortical regions.

Begin with tracers

Saffman-Turner hypothesis: Collisions should occur uniformly between any two regions that possess the same velocity gradient magnitude

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Discrepancy: Flow structures that are intrinsically more effective at causing collisions. Γ ^{IC}ARDO, AGASTHYA, AGASTHYA, AGASTHYA, AGASTHYA, AND RAYA, AGASTHYA, AGASTHYA, AGASTHYA, AGASTHYA, AGASTHYA, AGASTHYA, AGASTHYA, AGASTHY

Collisions *over sample* straining regions relative to where particles reside Comsions *Over sample* straining regions relative to writere particles reside

entardo, Agasthya, Govindarajan, and Ray, Phys. Rev. Fluids (Rapids) 4, 032601(R) (2019)

Take home messages FIG. 3. (a) Contours corresponding to *P*(*R*²

• Collisions *over sample* straining regions relative to where particles reside and theme moddaged
● Collisions *over sample s*training regions relative to where particles reside **Conditions doe regions dominated by strain and vorticity regions its various distributions of the conditions o**

Picardo, Agasthya, Govindarajan, and Ray, Phys. Rev. Fluids (Rapids) 4, 032601(R) (2019)

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- Collisions *head-on* or *rear-end* in straining regions; perpendicular in vortical regions indicate regions does regions dominated by strain and voltatively. Conditional probability distributions of the strain probability distribution of (c) the strain probability distribution of (c) the stribution of (c) the st complete the angle of the angle of the shall may region b, perpendicular in voltical regions of panel (a), and α

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- Lower approach velocities in vortical regions **b** Lower approach velocities in vortical regions **velocity is**

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why straining regions are more effective at generating regions are more effective at generating collisions. The
The collisions are more effective at generating collisions. The collisions are collisions are collisions at ge

In a given time interval, fewer particles will collide in vortical regions compared to straining regions

Picardo, Agasthya, Govindarajan, and Ray, Phys. Rev. Fluids (Rapids) 4, 032601(R) (2019)

How does particle inertia affect this picture?

Picardo, Agasthya, Govindarajan, and Ray, Phys. Rev. Fluids (Rapids) 4, 032601(R) (2019)

- Large Stokes: Particle dynamics decorrelated from the flow
- St ~ 0.3: Strongest effect for strain-dominated collisions Fig. 4. (a) Average Stones. Farticle uyitatifies decorrelated from the flow
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function of St, and normalized by the domain-average number density α **and velocity for** α **and** α

collisions in vortical and straining regions, normalized by *properties* • Preferential Concentration: Higher number of collisions

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- Preferential Concentration: Higher number of collisions collisions in vortical and straining regions, normalized by *properties*
	- **• Enhanced relative velocity** This surprising result is an outcome of the distinct flow topologies of the distinct flow top

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particles to approach each other differently. Figure 3(b) presents the distribution of the cosine of the Larger approach rates due to flow structures and not preferential concentration seems to be the the relative velocity vertor considers and the separation vector (*Xp* 2) and the separation vector (*Xp* 2) and the separation vector (*Xp* 2) and the separation of α primary driver for collisions

fraction of f and f the velocity difference between particles contributes to the f rate of α Perrin & Jonker, Phys. Rev. E 89, 033005 (2014) Perrin & Jonker, J. Fluid Mecjh. 792, 36 (2016)

Picardo, Agasthya, Govindarajan, and Ray, Phys. Rev. Fluids (Rapids) 4, 032601(R) (2019) Perrin & Jonker, Phys. Rev. E 89, 033005 (2014)
Picardo, Agasthya, Govindarajan, and Ray, Phys. Rev. Fluids (Rapids) 4, 032601(R) (2019)

Can vortical and straining regions conspire to enhance collisions?

Picardo, Agasthya, Govindarajan, and Ray, Phys. Rev. Fluids (Rapids) 4, 032601(R) (2019)

Geometry of the flow lead to particles in intense vortex tubes to be ejected into strong straining sheets leading to collisions

Picardo, Agasthya, Govindarajan, and Ray, Phys. Rev. Fluids (Rapids) 4, 032601(R) (2019) $C_{\rm c}$ larger fraction of the velocity gradient in straining z is translated into the velocity gradient into

Geometry of the flow lead to particles in intense vortex tubes to be ejected into strong straining sheets leading to collisions

FIG. 1. Representative somalise backward-in-time Ladran Use backward-in-time Lagrangian calculations to trace where colliding particles come from

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Particles that collide quickly originate in intense vortices and collide in regions of intense strain

Picardo, Agasthya, Govindarajan, and Ray, Phys. Rev. Fluids (Rapids) 4, 032601(R) (2019)

Geometry of the flow lead to particles in intense vortex tubes to be ejected into strong straining sheets leading to collisions

FIG. 1. Representative somalise backward-in-time Ladran Use backward-in-time Lagrangian calculations to trace where colliding particles come from

discrepancy, not simulated and concentration, but also by preference \mathbf{r} **Collisional velocities systematically larger when they structures conspire** PICARDO, AGASTHYA, GOVINDARAJAN, AND RAY **Particles that collide quickly originate in intense vortices and collide in regions of intense strain**

Summary

- Straining regions more effective at generating collisions: Larger proportion of *head-on* or *rear-end* collisions
- A larger fraction of the velocity gradient in straining zones is translated into the particle approach velocity.
- Intense vorticity and strain, cohabiting as vortex-strain worm-rolls, conspire to generate violent collisions.
- Preferential concentration, caustics and other mechanisms only a part of the story: Turbulence matters!

Conclusions & Perspective

Summary

- Straining regions more effective at generating collisions: Larger proportion of *head-on* or *rear-end* collisions
- A larger fraction of the velocity gradient in straining zones is translated into the particle approach velocity.
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PHYSICAL REVIEW E **99**, 063107 (2019)

Understanding droplet collisions through a model flow: Insights from a Burgers vortex

Lokahith Agasthya,^{1,2,3,*} Jason R. Picardo,^{2,†} S. Ravichandran,^{4,‡} Rama Govindarajan,^{2,§} and Samriddhi Sankar Ray^{2,¶} ¹*Indian Institute for Science Education and Research, Pune, 411008, India*

²*International Centre for Theoretical Sciences, Tata Institute of Fundamental Research, Bangalore 560089, India* ³*Department of Physics and INFN, University of Rome Tor Vergata,Via della Ricerca Scientifica 1, 00133 Rome, Italy*

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We investigate the role of intense vortical structures, similar to those in a turbulent flow, in enhancing collisions (and coalescences) which lead to the formation of large aggregates in particle-laden flows. By using a Burgers vortex model, we show, in particular, that vortex stretching significantly enhances sharp inhomogeneities in spatial particle densities, related to the rapid ejection of particles from intense vortices. Furthermore our work shows how such spatial clustering leads to an enhancement of collision rates and extreme statistics of collisional velocities. We also study the role of polydisperse suspensions in this enhancement. Our work uncovers an important principle, which, if valid for realistic turbulent flows, may be a factor in how small nuclei water droplets in warm clouds can aggregate to sizes large enough to trigger rain.

Picardo, Agasthya, Govindarajan, and Ray, Phys. Rev. Fluids (Rapids) 4, 032601(R) (2019)

Summary

- Straining regions more effective at generating collisions: Larger proportion of *head-on* or *rear-end* collisions
- A larger fraction of the velocity gradient in straining zones is translated into the particle approach velocity.
- Intense vorticity and strain, cohabiting as vortex-strain worm-rolls, conspire to generate violent collisions.
- Preferential concentration, caustics and other mechanisms only a part of the story: Turbulence matters!

Open Questions

- How are collisions affected when structures change?
- Modification of flow structures by new physical interactions: Examples include condensation, small-scale motion, elastic feedback from polymer additives.
- Role of gravitational settling